

# General Physics A: Final

## 1 Kepler's Laws of Planetary Motion

**Gravitational Binary System** The key goal of this problem is to derive Kepler's laws of planetary motion using the conservation laws we have learned this semester.

Consider two bodies located at  $\vec{x}_1$  and  $\vec{x}_2$ , with masses  $m_1$  and  $m_2$  respectively. Let us assume they interact solely through the force of gravity, without any external forces. Newton's second law now reads

$$m_1 \frac{d^2 \vec{x}_1}{dt^2} = - \frac{G_N m_1 m_2}{|\vec{x}_1 - \vec{x}_2|^2} \frac{\vec{x}_1 - \vec{x}_2}{|\vec{x}_1 - \vec{x}_2|}, \quad (1.0.1)$$

$$m_2 \frac{d^2 \vec{x}_2}{dt^2} = + \frac{G_N m_1 m_2}{|\vec{x}_1 - \vec{x}_2|^2} \frac{\vec{x}_1 - \vec{x}_2}{|\vec{x}_1 - \vec{x}_2|}. \quad (1.0.2)$$

**Galilean Transformation** To begin, explain why Newton's laws in equations (1.0.1) and (1.0.2) take the *same form* if we switch from the inertial frame  $\{\vec{x}_1, \vec{x}_2\}$  to another inertial frame  $\{\vec{x}'_1, \vec{x}'_2\}$  related by

$$\vec{x} = \vec{x}' + \vec{v}_0(t - t_0); \quad (1.0.3)$$

where  $\vec{v}_0$  and  $t_0$  are constants-in-time. The  $\vec{v}_0 \cdot t_0$  may be viewed as a constant displacement of the coordinate system; whereas  $\vec{v}_0 t$  relates two inertial frames that are moving at a constant velocity  $\vec{v}_0$  with respect to each other.

**Center-of-Mass Frame** Let  $\vec{X}_{\text{COM}}$  be the location of the center-of-mass of the above binary system; and let

$$\vec{\Delta} = \vec{x}_1 - \vec{x}_2. \quad (1.0.4)$$

Show that

$$\ddot{\vec{X}}_{\text{COM}} = 0, \quad (1.0.5)$$

$$\ddot{\vec{\Delta}} = - \frac{G_N M}{|\vec{\Delta}|^2} \frac{\vec{\Delta}}{|\vec{\Delta}|}; \quad (1.0.6)$$

where  $M = m_1 + m_2$  is the total mass; and each overdot denotes a time derivative.

**Energy and Angular Momentum Conservation** Explain why the following are constants-in-time:

$$E \equiv \frac{1}{2} \dot{\vec{\Delta}}^2 - \frac{G_N M}{|\vec{\Delta}|}, \quad (1.0.7)$$

$$\vec{L} \equiv \frac{1}{2} \vec{\Delta} \times \dot{\vec{\Delta}}. \quad (1.0.8)$$

**Kepler's 2nd Law** Explain why the conservation of eq. (1.0.8) implies the motion described by  $\vec{\Delta}$  must always lie on a fixed 2-dimensional (2D) plane. Next, explain why  $|\vec{L}| \equiv \dot{A}$  is in fact the area swept out by the  $\vec{\Delta}$ -motion per unit time. This is Kepler's 2nd law.

**Polar Coordinates** Since the binary system must move on a fixed plane, we may now use polar coordinates:

$$\vec{\Delta} = (r \cos \phi, r \sin \phi, 0). \quad (1.0.9)$$

Show that eq. (1.0.8) is equivalent to

$$\dot{\phi} = \frac{2\dot{A}}{r^2}; \quad (1.0.10)$$

and eq. (1.0.7) becomes

$$E = \frac{2\dot{A}^2 r'^2}{r^4} - \frac{G_N M}{r} + \frac{2\dot{A}^2}{r^2}. \quad (1.0.11)$$

Here, we have re-expressed  $r$  in terms of  $\phi$  and have denoted

$$r'(\phi) \equiv dr/d\phi. \quad (1.0.12)$$

**Kepler's 1st Law** Now, an ellipse  $(x, y)$  on the 2D plane obeys

$$\frac{x^2}{a^2} + \frac{y^2}{a^2(1-e^2)} = 1, \quad (1.0.13)$$

for some *semi-major* axis  $a > 0$  and *eccentricity*  $e$  subject to the constraint  $0 \leq e < 1$ . On the other hand, a hyperbola  $(x, y)$  on the 2D obeys

$$\frac{x^2}{a^2} - \frac{y^2}{a^2(e^2-1)} = 1, \quad (1.0.14)$$

but the eccentricity  $e$  is now subject to the constraint  $e > 1$ . It turns out that both equations (1.0.13) and (1.0.14) may be parametrized by

$$(x, y) = (ae, 0) + r(\phi)(\cos \phi, \sin \phi), \quad (1.0.15)$$

where  $(ae, 0)$  is the focus of the ellipse or hyperbola and

$$r(\phi) = \frac{a(1-e^2)}{1+e \cos \phi}. \quad (1.0.16)$$

(You *do not* need to derive this result.) Insert eq. (1.0.16) into eq. (1.0.11) and solve for  $a$  so that  $E$  is indeed a constant. With your solution for  $a$ , further demonstrate that

$$E = -\frac{1-e^2}{8} \left( \frac{G_N M}{\dot{A}} \right)^2. \quad (1.0.17)$$

Is  $E$  positive or negative for elliptical orbits? What about for hyperbolic orbits?

Notice you did not have to solve Newton's 2nd law in eq. (1.0.6)! But you have now shown that  $r(\phi)$  in fact describes either an ellipse or a hyperbola, provided  $a$  yields a constant  $E$ .<sup>1</sup> This is the (generalized) Kepler's 1st law.

**Kepler's 3rd Law** Finally, use the total area of the ellipse

$$A = \pi a^2 \sqrt{1 - e^2} \quad (1.0.18)$$

to deduce Kepler's 3rd law. For elliptical orbits, and with  $T$  denoting a period, show that

$$T^2 = \frac{4\pi^2 a^3}{G_N M}. \quad (1.0.19)$$

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<sup>1</sup>You are given the form of the solution in eq. (1.0.16); extra credit if you can *derive* it by integrating eq. (1.0.11) directly.